

Ising Lattice

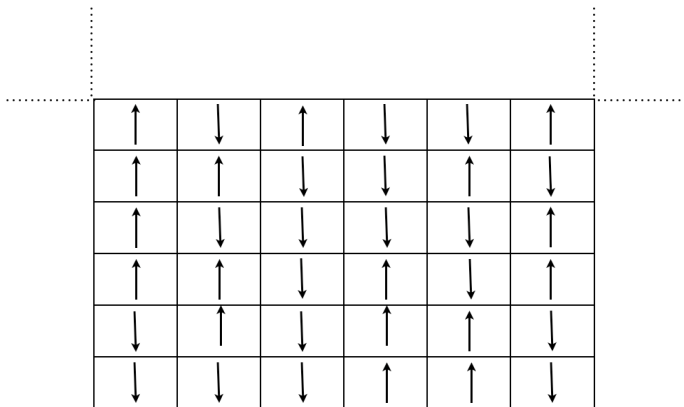
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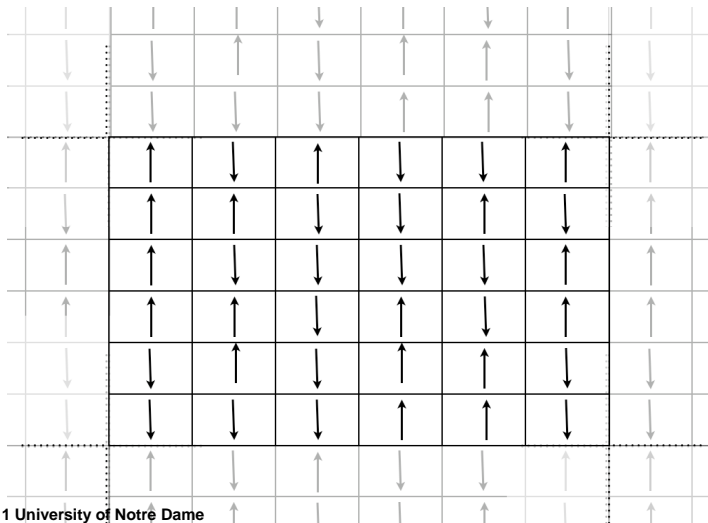
Ising Lattice Model

- Consider a 2-D lattice of “spins” up or down



Ising Lattice Model

- We replicate in all directions to add “periodic boundary conditions”



Ising Lattice Model

Consider a system of N spins on a lattice. In the presence of an external magnetic field, H , the energy of a particular state ν is

$$E_\nu = - \sum_{i=1}^N H s_i - J \sum_{ij} s_i s_j \quad (1)$$

- ▶ First term: energy due to individual spins coupling with external field
- ▶ Second term: energy due to interactions *between* spins.
- ▶ Assume that only nearest neighbors interact
- ▶ J is *coupling constant*, and describes the interaction energy between pairs.
- ▶ When $J > 0$, it is energetically favorable for neighboring pairs to be aligned.

Ising Lattice Model

If J is large enough (or temperature low enough), the tendency for neighboring spins to align will cause a cooperative phenomena called *spontaneous magnetization*.

- ▶ Physically: caused by interactions among nearest neighbors propagating throughout the system
- ▶ A given magnetic moment influences alignment of spins separated by a large distance
- ▶ Such long range correlations associated with long range order; lattice can have net magnetization *in the absence of external magnetic field*.
- ▶ Magnetization defined as

$$\langle M \rangle = \sum_{i=1}^N s_i \quad (2)$$

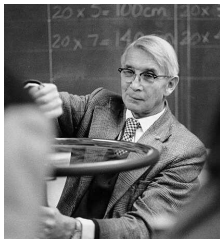
- ▶ A non-zero $\langle M \rangle$ when $H = 0$ is called *spontaneous magnetization*

Ising Lattice Model

- ▶ Temperature where system exhibits spontaneous magnetization is the *Curie temperature* (or critical temperature), T_c
- ▶ T_c is the highest temperature for which there can be a non-zero magnetization in the absence of an external magnetic field
- ▶ For $T_c > 0$, Ising model undergoes an *order-disorder transition*
- ▶ Similar to a phase transition in a fluid system - simple model of a fluid
- ▶ No order-disorder transition in 1-D only 2-D and 3-D

Ising Lattice Model

- ▶ 1-D Ising model solved analytically by Ernst Ising in his 1924 PhD thesis



- ▶ Historical note: German Jew who fled Europe, ending up in Peoria, IL as physics teacher at Bradley University
- ▶ Never published again after WWII
- ▶ Died in 1998 - see obituary:

<http://www.bradley.edu/las/phy/personnel/isingobit.html>

Ising Lattice Model

Lars Onsager showed in the 1940s that for $H = 0$, the partition function for a two-dimensional Ising Lattice is

$$Q(N, \beta, 0) = [2 \cosh(\beta J) e^I]^N \quad (3)$$

where

$$I = (2\pi)^{-1} \int_0^\pi d\phi \ln\left(\frac{1}{2}[1 + (1 - \kappa^2 \sin^2 \phi)^{1/2}]\right)$$

with

$$\kappa = 2 \sinh(2\beta J) / \cosh^2(2\beta J)$$

This result was one of the major achievements of modern statistical mechanics

Ising Lattice Model

- ▶ It can be shown that

$$T_c = 2.269J/k_B \quad (4)$$

- ▶ Furthermore, for $T < T_c$, the magnetization scales as

$$\frac{M}{N} \sim \alpha(T_c - T)^\lambda \quad (5)$$

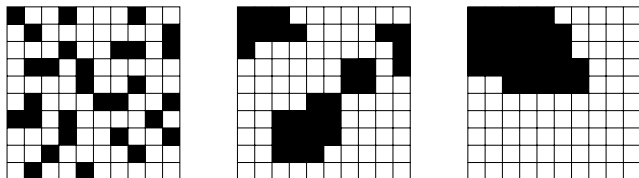
- ▶ Should be possible to perform Metropolis simulation of this!

Properties of Ising Lattice

- ▶ Physically, Ising lattice shows many of the characteristics of a fluid
- ▶ Magnetic susceptibility, $\chi = (\langle M^2 \rangle - \langle M \rangle^2)/k_B T$, diverges at critical point
- ▶ Local magnetization fluctuations become very large near critical point, similar to density fluctuations near critical point of a fluid
- ▶ Small variations in $k_B T/J$ lead to spontaneous phase changes

Properties of Ising Lattice

- ▶ The *correlation length* (distance over which local fluctuations are correlated) is unbounded at T_c



As T approaches T_c , correlation length increases

Figure: T_c is being approached from left to right. The correlation length of like regions (i.e. black and white squares) increases. At T_c , an order-disorder transition occurs, analogous to condensation

Ising Lattice Algorithm

How to simulate the 2-D Ising lattice?

1. Choose an initial state of spins (it will not matter)
2. Choose a site i
3. Calculate the change in energy ΔE if the state if i is changed
4. If energy is lowered by change, accept the change. If not...
5. Generate random number $0 < \zeta < 1$
6. If $\zeta < \exp(-\beta\Delta E)$ accept the change. Otherwise, don't
7. Repeat

Compute $\langle M \rangle$, $\langle E \rangle$ and fluctuations versus T . Let's do it!

Ising interactive Algorithm

Take a look at the source code `ising2d.f`

- ▶ Lattice array or 0 or 1:

```
latt(ix,iy)
```

- ▶ Largest allowable lattice size: 40 X 40
- ▶ Initialize it in one of four choices
- ▶ Compute the initial energy

```
CALL echeck
```

Ising interactive Algorithm

Core Metropolis algorithm is in routine “update”

```
do istep=1,nstep
  CALL update

ccc  sum E, M and fluctuations
      emean=emean+etot
      efluc=efluc+etot*etot
      rmmean=rmmean+dfloat(mtot)
      rmfluc=rmfluc+dfloat(mtot)*dfloat(mtot)
enddo                                ! istep = 1,nstep
```

Ising interactive Algorithm

Metropolis implementation for an nsize X nsize lattice

```
size=dfloat(nsize)
nsize1=nsize-1
```

```
ccc Pick a spin at random. Each call to
ccc update performs nsize^2 attempts
```

```
do iflip=1,nsize*nsize
ccc  (ix,iy) is a random position on lattice
    ix=int(size*uni())+1
    iy=int(size*uni())+1
```

```
ccc nhsum is a counter that sums the value
ccc of all spins adjacent to (ix,iy)
    nhsum=0
```

```
ccc    get positions adjacent to (ix,iy)
ccc    (remember PBC)
      do nhbr=1,4
        nhx=mod(ix+nhxl(nhbr)+nsizel,nsizel)+1
        nhy=mod(iy+nhyl(nhbr)+nsizel,nsizel)+1
        nhsum=nhsum+latt(nhx,nhy)
      enddo

ccc    compute energy felt by spin at (ix,iy)
      itest=nhsum*latt(ix,iy)
```


Here is the Metropolis algorithm

```
if (itest.le.0) then
ccc Automatically accept the move
ccc Sum total energy and magnetization
    mtot=mtot-2*latt(ix,iy)
    latt(ix,iy)=-latt(ix,iy)
else
ccc Select a random number on (0,1)
ccc Conditional accept (embe is precalculated E)
    if (test.le.embe(itest)) then
        etot=etot+dfloat(itest)*rj2
        mtot=mtot-2*latt(ix,iy)
        latt(ix,iy)=-latt(ix,iy)
    end if
ccc Reject the move. Keep everything the same
end if
```

Ising interactive

- ▶ Run Ising interactive at $T > T_c$. Note the value of the magnetization , fluctuations
- ▶ Dump the configuration and look at it. What do you see?
- ▶ Slowly reduce T and approach T_c . What happens?
- ▶ Start at low T ($T \approx 0.4$). What happens?
- ▶ Slowly raise T to approach $T_c = 2.26$ What happens?

Ising batch input file

```

ising.40.10K.log : name of log file
123456           : rdm nbr seed
mvst.1           : magnetization vs T filename
75              : unit (don't change)
evst.1           : energyvs T filename
76              : unit (don't change)
mflucvst.1       : magnetiz fluc vs T filename
77              : unit (don't change)
eflucvst.1       : energy fluc vs T filename
78              : unit (don't change)
1               : 1=random,2-inter,3=check,4=read
5000            : number of equilibration steps
2              : n different runs
40 10000 4       : nsize nstep kT/J for run
40 10000 2       : nsize nstep kT/J

```